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To



RESEARCH MEMORANDUM

EFFECTS OF AIRFOIL PROFILE ON THE TWO-DIMENSIONAL

FLUTTER DERIVATIVES FOR WINGS OSCILLATING

IN PITCH AT HIGH SUBSONIC SPEEDS

By John A. Wyss and James C. Monfort

Ames Aeronautical Laboratory
Moffett Field, Calif.

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NATIONAL ADVISORY COMMITTEE FOR AERONAUTICS

WASHINGTON

May 24, 1954





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RESEARCH MEMORANDUM

EFFECTS OF AIRFOIL PROFILE ON THE TWO-DIMENSIONAL FLUTTER DERIVATIVES FOR WINGS OSCILLATING

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SUMMARY

Aerodynamic lift and moment flutter derivatives were determined at high subsonic speeds for a series of two-dimensional airfoils varying in thickness and thickness distribution. The wings were sinusoidally oscillated about the quarter-chord axis at Mach numbers from about 0.5 to 0.9. The corresponding reduced frequency ranges varied from 0.045 to 0.45 at M = 0.5 and from 0.025 to 0.25 at M = 0.9. An evaluation of the results indicated that wing profile and angle of attack have major effects on the flutter derivatives at speeds exceeding the Mach number for steady-state lift divergence. In general, at supercritical Mach numbers the trends of the magnitudes of the oscillatory lift coefficients were qualitatively indicated by the trends of the nonoscillatory coefficients, with phase angles, except for the 12-percent-thick airfoil, having only moderate deviation from subsonic theory. The variations in the magnitude of the moment derivative and in its phase angle, resulted in a trend toward instability at supercritical Mach numbers. In particular, for airfoils of equal thickness the effect of an extreme forward location of maximum thickness was destabilizing in that negative -aerodynamic damping existed, implying the possibility of a single degree of freedom type of flutter. Decreasing airfoil thickness delayed the large deviation from subsonic theory to higher Mach numbers.

INTRODUCTION

This report is concerned with the evaluation of the effects of airfoil profile on the lift and moment flutter derivatives as measured, by means of pressure cells, on harmonically vibrating two-dimensional wings at high subsonic speeds. It is well-known that theory does not account properly for such factors as flow separation and shock formation, hence, the aircraft designer must of necessity look to experimental values

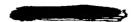


whenever such mixed-flow conditions are encountered. Numerous previous investigations at lower speeds, such as those by Clevenson and Widmayer (ref. 1) and by Halfman (ref. 2), may be cited. With the use of a different measuring technique, the present work extends these previous investigations to higher Mach numbers so that emphasis may be placed upon supercritical speeds for which information is meager or nonexistent.

Since wing profile may be expected to have a significant effect on mixed-flow conditions, several models were used to determine the effects of wing thickness and thickness distribution on the flutter derivatives. NACA 65A series symmetrical airfoils, 12, 8, and 4 percent thick, were used along with two other 8-percent-thick airfoils with their maximum thickness at about 16 and 63 percent of the wing chord. The models were oscillated about the quarter-chord axis at Mach numbers from 0.5 to 0.9 with reduced frequency ranges from 0.045 to 0.45 and from 0.025 to 0.25, respectively. Reynolds numbers, based on the airfoil chord, varied from 5 to 8 million.

SYMBOLS

a	velocity of sound in undisturbed air, ft/sec
ъ ;	wing semichord, ft
c _l	dynamic section lift coefficient
c _m	dynamic section moment coefficient about quarter point of chord
f	frequency of oscillation, cps
k	reduced frequency, $\frac{\omega b}{V}$
М .	Mach number, $\frac{V}{a}$
Ma	oscillatory aerodynamic section moment on wing about axis of rotation, positive with leading edge up
P_{α}	oscillatory aerodynamic section lift on wing, positive upwards
đ	free-stream dynamic pressure, 1b/sq ft
Λ .	free-stream velocity, ft/sec



α	oscillatory angular displacement (pitch) about axis of rotation, positive with leading edge up, radians
$\alpha_{\underline{m}}$	mean angle of attack about which oscillation takes place, deg
θ	phase angle between oscillatory moment and position $\alpha,$ positive for moment leading $\alpha,$ deg
φ	phase angle between oscillatory lift and position α , positive for lift leading α , deg
(i)	circular frequency, 2πf, radians/sec
de _l	magnitude of dynamic lift-curve slope, $\frac{P_{\alpha}e^{-i\phi}}{2bq\alpha}$, per radian
dc _m	magnitude of dynamic moment-curve slope, $\frac{M_{\alpha}e^{-i\theta}}{4b^2q\alpha}$, per radian
$\left \frac{\mathrm{dc_m}}{\mathrm{d\alpha}} \right \sin \theta$	aerodynamic damping component in phase with angular velocity

APPARATUS AND METHOD

Models and Instrumentation

The 12- and 8-percent-thick airfoils, NACA 65A012, 65A008, 2-008, and 877A008 profiles, were of wood-rib and wood-stressed-skin construction built around steel spars at the quarter chord, which was the axis of rotation. Several wood spars at other chordwise locations were used to minimize spanwise twisting since the models were driven from one side. The 4-percent-thick model, of NACA 65A004 profile, was machined from solid aluminum with a parting line in the chord plane. The upper and lower halves of this model were bolted and doweled together. Each model had a chord of 24 inches and a span of 18-1/4 inches. The gaps between the ends of the models and tunnel walls were sealed with sliding spring-loaded felt pads or brass strips which moved with the models.

¹An NACA 847AllO airfoil was modified to a symmetrical section by using the lower surface coordinates for both upper and lower surfaces and then reducing the thickness ratio to 8 percent.



In figure 1, the model profiles are illustrated to show the variation of thickness and thickness distribution. The reference model, NACA 65A008, is marked to indicate the locations of the pressure cells. Model instrumentation consisted of 15 flush-type pressure cells (see refs. 3 and 4) and 15 pressure orifices along the midspan of each surface of each model. The pressure orifices adjacent to each pressure cell were used for two purposes: (1) as a means to determine the time-average chordwise pressure distribution with the use of a multiple mercury manometer, and (2) to provide an internal reference pressure for the pressure cells. The tubes from each cell and from the adjacent pressure orifice were interconnected at the manometer. In order that the internal reference pressure of the pressure cells would be essentially steady, about 50 feet of 1/16-inch tubing was used from the orifice to the manometer and back to the pressure cell.

Two 14-channel oscillographs were used to record the instantaneous electrical difference of the output of each pair of cells (proportional to the pressure difference between the upper and lower surface at each chord station) and to record the summation of all cells (proportional to the variation of the lift force). The output of an NACA slide-wire position transducer, proportional to the model angle of attack, was simultaneously recorded.

Tunnel, Model Drive System, and Tests

The models were oscillated in the two-dimensional test section in the Ames 16-foot high-speed wind tunnel (ref. 5). The two-dimensional channel was about 20 feet long and 16 feet high. A view of a model in place and a diagrammatic sketch of the drive system are presented in figure 2. The drive rods and sector arm attached to the model were contained within one of the channel walls.

Records were obtained with Mach number and mean angle of attack constant for frequencies from 4 to 40 cycles per second at intervals of 4 cycles per second and for an amplitude of ±1°. Data are presented for mean angles of attack of 0° and 2° and for Mach numbers from 0.5 to about 0.9. Sample oscillograph records which illustrate the necessity for harmonic analysis at the higher Mach numbers are given in figure 3. The lift was evaluated by a 12-point harmonic analysis of three consecutive cycles of the sum trace. The pitching moment was evaluated by a 12-point harmonic analysis of the individual cell traces for one cycle.

Since the investigation was conducted in a closed-throat tunnel, the effects of wind-tunnel resonance must be accounted for either by avoiding conditions in which tunnel-wall effects are significant or by correcting the results for the effects of the tunnel walls (refs. 6 and 7). Calculations made at the Langley and Ames Laboratories employing

the single-doublet-line, single-control-point solution described in reference 7 yielded the following results for a tunnel height of 16 feet, wing chord of 2 feet, and Mach number of 0.7: At frequencies of 10, 20, and 40.66 cycles per second, the magnitudes of the coefficients were increased by 3.8, 5.0, and 4.7 percent, respectively, due to the presence of the tunnel walls. These results indicate that, for the conditions of the calculations, the effect of the tunnel walls was small. However, for mixed-flow conditions, the application of such corrections based on potential flow would be questionable; hence, to minimize tunnel-wall effects, all data obtained at frequencies within 10 percent of the tunnel resonant frequency (refs. 6 and 7) have been omitted. Although the use of such a procedure does not mean tunnel-wall effects have been completely eliminated over the entire frequency range, it is felt that tunnel-wall effects are not a predominant factor in the trends of the data.

For a discussion of other factors influencing the precision of the data, the reader is referred to references 3 and 4.

RESULTS AND DISCUSSION

A tabulation of the measured derivatives is contained in tables I, II, III, IV, and V for the NACA 65A012, 65A008, 65A004, 2-008, and 877A008 airfoils, respectively. The results concerning lift derivatives are first discussed and are presented in figures 4 to 10, followed by a discussion and the presentation of the moment derivatives in figures 11 to 15.

Lift

Experimental values for the reference model for three representative Mach numbers are presented in figure 4 as a function of reduced frequency. In this figure, as in subsequent figures, the absolute magnitude of the flutter derivative is expressed in terms of the slope of the lift curve per radian and the corresponding phase-angle relationship between the lift vector and model angle of attack in degrees. Theoretical values at Mach numbers of 0.5, 0.6, and 0.7 may be obtained from the work of Dietze (refs. 8 and 9), and at Mach numbers of 0.8 and 1.0 from Minhinnick (ref. 10) and Nelson and Berman (ref. 11), respectively.

In this figure it may be noted that at 0.49 and 0.79 Mach numbers the flutter derivatives tend to increase with increasing reduced frequency; furthermore, there seems to be a large variation in the phase angle at low values of reduced frequency at 0.79 Mach number. However,



Mach number appears to have had a greater effect on the data than did frequency at 0.91 Mach number.

Typical results as a function of Mach number are presented in figure 5 for the reference model, the NACA 65A008 airfoil. The lines showing the theoretical values are identified at one end by the frequency in cycles per second to which they pertain. Since theoretical values have been computed in the cited references only at certain Mach numbers which have already been indicated, an interpolation was necessary to obtain values at intermediate Mach numbers. Although such an interpolation inherently involves some error, a consistent set of values was nevertheless established and was used for the purpose of determining the effects of varying airfoil shape.

To distinguish between the various frequencies, the experimental and theoretical values are each faired with the same type of line. For example, the experimental and theoretical values for a frequency of 8 cycles per second are each shown with a solid line. Examination of the experimental data for a frequency of 8 cycles per second indicates that the trends of both experiment and theory were the same at low Mach numbers. As Mach number increased, a large decrease in the magnitude of the experimental derivative occurred, accompanied by a variation of phase angle such that the trend toward increasing lag was reversed. Although the agreement with theory was not precise at the lower Mach numbers, it may be seen that the general trends for all frequencies were nearly the same.

The data from figure 5 are presented in a different form in figure 6; the experimental magnitude has been divided by the theoretical magnitude, and the theoretical phase angle has been subtracted from the experimental phase angle. These quantities are also shown as a function of Mach number. If the experimental and theoretical values exactly agreed, the ratio of the magnitudes of the derivatives would be 1, while the difference in phase angle would be 0. The faired lines represent the average deviation from theory for the entire frequency range up to 40 cycles per second.

It is of interest to note that the individual points do not indicate an entirely random scatter about the mean line for the various frequencies. For example, examination of the points for 40 cycles per second in the top portion of the figure shows that these points are usually the uppermost value at each Mach number. Hence, this figure not only provides some indication of the range of the experimental values, but illustrates the fact that, although the values depend on frequency, the general variations with Mach number are represented by the faired average curves.

The use of the average deviation from theory appears to be justified since it is representative of each model. For example, in figure 6 it may be noted that all the experimental points lie within a comparatively



narrow band along the faired curves with the exception of the higher frequencies in the upper portion of the figure. In fact, a band of width ±0.15 in the upper portion of the figure and a band of width ±10° in the lower portion of the figure would contain about 80 percent of all the experimental points. These results are typical of all the models. It might be noted that the averaging process used has the effect of removing frequency as a parameter. It should be noted that each model was oscillated at the same amplitude and through the same range of frequencies, hence the average deviation from theory indicates the over-all effects of airfoil shape and the general trends of the data.

Effect of thickness distribution. The effects of the variation of thickness distribution as indicated by the curves showing the average deviation from theory over the frequency range tested are summarized in figure 7 for mean angles of attack of 0° and 2°. It would appear from this figure that the main effect of the chordwise location of maximum thickness was on the magnitudes of the derivatives rather than on phase angles, although no systematic trend is apparent.

Effect of wing thickness.— The results showing the effects of wing thickness are presented in figure 8. At an angle of attack of 0° , wing thickness appears to have had a much more pronounced effect than wingthickness distribution (fig. 8(a) as compared to fig. 7(a)). As might be expected, the primary effect of reducing wing thickness was to delay any large deviation from theory to a higher Mach number.

At an angle of attack of 2° (fig. 8(b)), large differences over the entire range of Mach numbers occurred between the models in the magnitudes of the derivatives.

Comparison with steady-state results. In order to examine whether any relation existed between unsteady and steady-state results, a comparison with steady-state results obtained from the time-average chord-wise pressure distributions for mean angles of attack of 0° and 2° is made in figures 9 and 10. In these figures, the steady-state data have been normalized with the Prandtl-Glauert value of the theoretical lift-curve slope. It may be recalled that the Prandtl-Glauert curve is also obtained as an end condition as the frequency of oscillation approaches zero.

Examination of these figures indicates that although there appears to be some parallelism or similarity between the steady and unsteady curves, the comparison between the steady and unsteady values is at best only qualitative. For example, in neither figure 9 nor figure 10 do the unsteady and steady-state curves coincide throughout the entire range of Mach numbers. It should also be noted that, with the exception of the NACA 65A012 airfoil at a mean angle of attack of 2° (fig. 10(b)), the unsteady values approached theory more closely than did the steady-state



values, particularly at the lower Mach numbers, that is, from M=0.5 to 0.7. Although the effect of the higher frequencies in increasing the level of the curves for the unsteady case may in part account for the differences between the curves, this effect is small. However, the one characteristic that is common to both the unsteady and steady curves in almost every case is a trend toward a reduction in magnitude at the highest Mach numbers. The Mach number at which this trend initiates cannot be precisely delimited, nevertheless, for the three NACA 65A-series airfoils at a mean angle of attack of 0° (fig. 10(a)), the unsteady lift trend appears to be associated with the steady-state flow changes which occur above the Mach number for lift divergence.

It would therefore appear that as a first approximation the Mach number for lift divergence may be taken as a criterion for the onset of significant changes in the trends of the unsteady values, and that this trend toward a decrease in the magnitude of the unsteady values is related to the trend of the steady-state data. It should be pointed out that this conclusion is not as evident for the NACA 2-008 and 877A008 airfoils (fig. 9) and for the NACA 65A004 airfoil at a mean angle of attack of 2° (fig. 10(b)), since these figures indicate that the correlation between the Mach number for lift divergence and the initiation of a downward trend of the unsteady values is not precise and they may differ by as much as 0.1. However, it is felt that there is sufficient evidence presented in figures 9 and 10 to indicate that steady-state values may prove useful as a qualitative indication of the trends of the unsteady-state coefficients at supercritical Mach numbers.

For the steady-state condition the phase angle is, of course, zero; therefore no corollary for the phase angle with relation to the oscillatory condition is possible. However, except for the 12-percent-thick wing, the phase angle shows only a moderate deviation from theory throughout the speed range of the present investigation.

Moment

The moment derivatives for the reference model as a function of reduced frequency for several Mach numbers are presented in figure 11 and as a function of Mach number in figure 12. A comparison of these figures indicates that even though there may have been a greater effect due to frequency on the moment derivatives than had been the case for the lift derivatives, from figure 12 it appears that the effects of Mach number are similar for all frequencies. Hence, the effects of airfoil profile are again compared on the basis of the faired average curves in figure 12 which represent the average deviation from theory over the entire frequency range.



In contrast to the lift results previously presented in figure 6, the magnitudes of the moment derivatives greatly exceeded the theoretical values, along with a much larger variation of phase angle as compared with theory. These results may be attributed to the fact that the comparison is between very small quantities in regard to the magnitude of the derivatives, since the moment is taken about the quarter-chord axis, and to small movements of the center of pressure which would be reflected in large changes of phase angle. The general trends of the results, nevertheless, are represented by the faired average curves.

Effect of thickness distribution. The effects of the variation of the chordwise location of maximum thickness are shown in figure 13. An apparent characteristic of the NACA 2-008 airfoil, with a forward location of maximum thickness, is a large shift toward a lagging phase angle as Mach number increased above 0.8, such that the phase angle lagged theory by 80° and 90° at angles of attack of 0° and 2°, respectively. The effects of such large shifts in phase angle are discussed in relation to subsequent figures.

Effect of wing thickness. The effects of wing thickness on the moment derivatives are shown in figure 14. As might be expected, the primary effect of decreasing wing thickness was again to delay any large variations to a higher Mach number.

Instability. Since there was such a large variation at the higher Mach numbers from the subsonic theoretical values, it is of basic importance to examine the damping-moment derivatives directly to determine whether instability, or the existence of negative aerodynamic damping (implying the possibility of a single degree of freedom type of flutter), which is not predicted by the theory, existed at these speeds. The average damping-moment derivatives for the entire frequency range are therefore presented in figure 15. Also included in this figure are dashed lines indicating average values derived from theory for the corresponding frequency range.

The effect of wing-thickness distribution on aerodynamic damping is shown in figure 15(a) for each mean angle of attack. It may be noted that there was a trend toward instability for each model, with the NACA 2-008 airfoil becoming abruptly unstable at about 0.85 Mach number at 0° and 2° angles of attack. It would appear that stability about the quarter-chord axis increased as maximum thickness was moved toward the trailing edge.

The effect of wing thickness on the aerodynamic damping moment is shown in figure 15(b) for each angle of attack. Although the trend toward instability does not appear at 0° angle of attack for the NACA 65AOO4 profile, the susceptibility of the thinner wing to negative aerodynamic damping is clearly indicated at the 2° mean angle of attack.





CONCLUSIONS

Within the limitations of speed range and angle-of-attack variation of the investigation, the following general conclusions may be drawn:

- 1. Section profile has a major effect on the flutter derivatives at speeds exceeding the Mach number for steady-state lift divergence.
- 2. It appears that the variation in angle of attack has an effect as important as the effect of the variation in profile.
- 3. In general, at supercritical Mach numbers, a qualitative evaluation of the results indicated that the trends of the magnitudes of the oscillatory lift coefficients were indicated by the trends of the non-oscillatory lift coefficients, with phase angles, except for the 12-percent-thick model, showing only a moderate deviation from theory.
- 4. The variations in the magnitude of the moment derivative and in its phase angle, resulted in a trend toward instability at supercritical Mach numbers. In particular, for airfoils of equal thickness the effect of an extreme forward location of maximum thickness was destabilizing in that negative aerodynamic damping existed, implying the possibility of a single degree of freedom type of flutter.

Ames Aeronautical Laboratory
National Advisory Committee for Aeronautics
Moffett Field, Calif., Mar. 24, 1954

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TABLE I.- MEASURED FLUTTER DERIVATIVES FOR THE NACA 65A012 AIRFOIL

			$\alpha_{m} = 0^{\circ}$				α _m = 2°						
к	k	tu	de	φ	dc _m	8	м	k	ω	dc 1	φ	dc _m	в
0.491	0.103 .184 .282	57:0 101.8 155.9	6.394 5.466 5.099	351.8 358.8 355.5	===		0.491	0.058 .094 .136 .187	31.7 51.1 74.1	6.520 5.578 5.574 4.989	354.6 354.1 0.0	= = =	:
.590	.077 .152 .229	51.6 101.3 153.2	7.083 6.056 5.319	351.7 351.9 351.2	===			.238 .287 .328	102.5 130.1 157.1 179.5	4.987 5.341 5.058	5.3 4.5 0.0 12.4		
.633	.074 .111 .144 .183 .218 .252 .320	52.6 79.2 103.0 130.9 155.9 180.0 228.5 256.5	5.745 5.068 5.299 4.661 4.449 4.036 3.913 4.259	355.0 355.5 355.5 357.3 358.0 349.6 15.0	0.531 .590 .585 1.008	342.1 317.6 305.0 311.6	.590	.469 .048 .076 .120 .152 .198 .233 .347 .384	256.5 31.5 50.5 79.4 100.8 131.2 154.0 229.3 254.4	4.823 6.523 6.262 5.925 5.965 5.488 5.213 4.744 5.426	29.4 352.4 345.5 347.7 354.2 352.5 347.7 9.2 16.1	0.559 .771 .739	3 ⁴ 1.
.682	.064 .097 .130 .163 .197 .264 .293	49.8 76.0 101.6 127.3 153.7 206.2 229.3 254.4	7.918 7.332 6.855 5.533 5.765 4.362 5.118 4.932	344.4 339.7 348.2 337.3 346.4 2.5 0.8 0.4	.595 .658 .745 .554 .868	325.4 310.6 279.7 291.5 278.8	.682	.044 .066 .101 .131 .163 .196	34.6 52.0 80.2 103.7 128.9 154.8 232.7	6.216 5.833 5.506 5.224 5.055 4.528 4.290	354.5 349.2 347.3 0.4 348.8 342.7 15.3		
.731	.062 .098 .121 .156 .247 .280 .308	51.2 81.2 100.6 129.4 205.6 232.7 256.1	8.080 8.454 7.092 6.092 5.187 5.299 5.018	348.1 339.5 339.5 328.9 356.2 355.2 4.4	.634 .675 .647	326.9 304.6 283.0 279.5	.731	.321 .041 .060 .093 .122 .153 .240	253.7 34.6 51.2 79.2 104.5 130.4 204.9	4.329 6.788 6.050 5.566 5.437 5.280 4.182	2.0 351.2 348.9 349.8 351.9 346.3 359.8	.698 .642 .721	340. 333. 315. 300.
.790	.057 .086 .114 .142 .199 .226 .256	52.2 77.8 103.9 129.3 180.9 205.3 232.7	8.576 8.362 7.476 6.137 4.771 4.588 5.285	343.5 337.9 336.4 327.1 351.7 348.9 356.6	.045	305.4 276.3 263.1	.790	.271 .299 .034 .056 .086 .115	231.3 255.4 30.9 50.8 77.6 103.5 125.1	4.375 4.282 6.377 5.981 7.353 6.628 5.099	0.6 358.2 353.8 347.9 343.3 341.9 333.5	1.084 .597 .606	340. 316.
.837	.052 .077 .104 .182	50.8 74.7 101.5 177.3	4.894 4.590 4.780 3.515	354.0 342.7 351.6 2.6	.612	301.6		.198 .225 .254 .279	178.5 202.7 228.8 251.6	3.861 4.047 4.196 4.895	353.3 348.9 358.7 0.9	.557 1.126	287.
	.207 .235 .262	200.9 228.4 255.4	3.597 4.444 5.123	12.2 16.5 359.9	.857 1.736	25 6. 9 222.0	.837	.031 .054 .080	29.6 51.8 77.7 100.2	4.318 4.580 4.775 4.570	355.4 357.6 353.3 356.9	.285 .606	340.8 306.
.885	.030 .049 .080 .097	30.7 50.3 82.1 99.7 153.2	.965 .641 1.725 1.884 2.681	47.0 92.7 59.9 47.4 41.2	2.719 3.117 2.436 1.939	348.1 348.0 311.4 314.4		.181 .208 .244 .261	175.5 201.7 236.5 252.3	3.654 4.068 4.601 5.379	348.1 351.4 13.9 5.1	675 1.316	259.2
	.176 .201 .223 .246	181.3 207.3 230.4 253.3	2.015 1.454 2.733 2.681	29.0 33.2 32.1 2.4	1.223	304.0	.885	.030 .049 .080 .097	30.7 50.3 82.1 99.7 153.2	3.497 2.751 3.032 2.403 2.647	347.1 345.4 359.2 342.9 353.5	1.256 .879 .712 1.389	356.8 350.2 340.7
								.177 .202 .225 .247	181.3 207.3 230.4 253.3	3.564 2.122 3.084 3.944	351.6 336.5 345.6 356.4	1.043	333.5





TABLE II.- MEASURED FLUTTER DERIVATIVES FOR THE NACA 65A008 AIRFOIL

			c _m = 0 ^c							α ²⁴ = 5 _c			
И	k	۵	de	φ	dCa da	0	¥	k	a 0	da _z	9	de	8
0.491	0.089	48.9	6.186	354-7			0.492		30.7	5.356	352.1		
	.184	78.2	5.638	348.3 353.3				.089	48.9 76.7	5.074 4.613	354.4		
	.234	128.5	5.319	357.0	1222		1	186	101.8	4.571	357.1		1:::/
	•280	153.9	6.250	1.9			l	.231	126.7	4.205	358.0		
	-322 -157	251.3	5.518	352.5 19.2		1:::	ĺ	.281	154.4	4.510 4.443	4.6 356.0	= = =	= =
			-	1			1	.463	254.4	4.828	24.2		
.590	.074	17.0	5.841	351.0	0.445	315.8	.590	.031	20.3	6.214	340 6	0.581	341.1
	.152	101.2	5.756	343.6	.626	312.2	کور. ا	.076	50.6	5.803	349.6	-557	331.6
	-191	127.2	5.854	347-5			l l	.117	77.8	5.263	346.4		
	.234	155.9	5.673	349.1	-775	289.8	i	.153 .189	102.0	5.237 4.863	348.0	.588	307.1
	.270 .347	179.5 231.0	5.612	6.7			ļ	-231	153.2	4.975 4.646	346.9	.701	284.4
	.383	254.4	6.787	9-5	1.213	290.0	İ	.271	180.0	4.646	345.5		
.680	.065	50.1	6.815	348.9			1	.346 .378	230.1	4.122 5.246	13.9	.964	281.3
	.102	79.4	6.312	347.0				!			1	l '	
- 1	.130 .167	101.0	6.062	338.4			.680	.038	29.3	6.618	351.7 346.8	.778	341.0
- 1	.199	154.4	5.652	338.2			1	.099	76.3	5.972	3+3-3		331.4
}	-295	228.8	6.067	2.0				•133	102.6	5-753	341.5	.811	309.7
	-329	255.1	6.389	2.3			-	.200	127.9	5.611 5.339	335.9 330.3	.871	282.0
-728	.060	50.4	7-392	340.9				.296	228.8	5,302	357.5		
	.092	77.2	7.005	339-5				-327	252.6	5.662	350.6	1.097	278.9
	.123 .154	102.5	6.535 6.028	335.5 331.0			.728	.037	30.7	7.311	350.6	.888	350.5
)	245	204.9 230.4	5.696	347.0			0,20	.053	22.8	7.347 6.759	343-1	.944	330.8
	.276 .305	230.4	6.297	348.9 352.4				.095	80.4	6.159	337.6 338.7		
- 1	•305	254.9	0.190	352.4	!		' '	156	131.3	6.247	1326.4	-972	304.7
.786	.058	2.3	7.999	339.5	.837	323.9		.245	206.2	5.508	349.3	-979	292.6
1	.086	78.7	7.381 6.851	336.1 326.9	.805	300.5		.277 .303	233.0	6.156 5.911	347.3	1.233	217.5
į	.143	130.2	6.132	320.9					1 1			1.233	211.0
1	.199 .225	181.4	5.394 5.525	348.3	.829	284.5	.761	.036	31.8	7.863	345.2		
1	.252	229.9	5.800	347.7	.029		1	.059 .091	51.8 80.0	6.883	341.9	===	===
	.279	254.4	6.848	345.7	1.367	271.4		.117	102.9	6.377	331.2		
.833	.050	47.7	7.488	335.9	. 4 54	309.4		.146	128.9 182.0	4.217	320.8 351.8		
	.080	76.1	6.865	332.5				.234	206.0	5.224	350.2		
	.105 .214	201.9	6.343 4.705	325.5 356.2	.487 .651	263.8	1	.261	229.8 255.5	5.879	346.7 354.4		
- 1	.238	227.6	5.124	353.5		291.9		.290	اردروعا) :		
Ì	.263	251.3	6.365	345.7	1.276	272.6	.786	.034	31.1	9.588	347.1	1.086	337-5
.879	.026	26.7	9.006	336.4	.223	198.6		.057	72.5 76.7	8.362 7.931	337.1 333.0	1.092	320.2
	.048	49.2	6,862	349.4	.203	281.9	!!	.114	104.1	7.931 7.520	326.8	1.053	286.4
	.074 .149	75.7 152.9	T-193	327.6	.093			-144	131.6	6.300	318.2 349.1		
- }	.174	178.8	4.299 4.315	347.6 353.4		155.9		.199	181.8	4.694 5.278	340.5	-973	275.8
Ì	.200	205.6	315 3.360	358.8	.517	245.9	1	.254	232.1	5.627	343.6		
- 1	.223 .248	229.0 255.1	5.990 6.383	348.1 337.5	1.157	189.3		.283	258.2	6.853	343.5	1.526	261.3
	1						.833	.030	29.4	8.034	347.4	.487	341.9
-917	.029	31.5 51.3	3.742 3.360	349.8 344.3	1.232	327.8		.055	53.5 79.4	7.464	338.4	-532	302.9
1	.073	_ 79.0	2.973	341_9	1.729	319.3	[[.107	104.4	6.595	325.5	.479	254.9
Ī	.118	127.7	3.437	346.7 348.4	: :::]]	.160	155.6	4.320	344.1	.206	284.9
- 1	.145 .167	156.3	3.654 3.634	348.4	1.161	300.6		.186	181.1 208.3	5.163 5.302	348.2 355.2	.678	286.3
l	.188]	180 3 202 4	3-515	346.1	.922	271.6) i	.236	230.7	6.237	355.2 354.0		
1	.214	230.5	4.228 3.997	349.5	1.440	252.4		.262	256.1	5.502	340.8	1.013	254.1
1	.250	د.سرے	コ・ファ	7-1-07	4.770	€.F	.879	.031	31.7	7.030	347.9	.380	311.2
l	1	1		1	f		"	-053	55.2	7.460	331.6	.29k	284.4
- 1	ı	1			- 1			.075	77.7	6.721 5.835	333.8 330.1	-323	198.4
1		1	1	1	1) [.150	155.9	4.641	336.5	.197	319.9
- 1	Ţ	ŧ	ł	ļ	[.174	180.7	4.870	351.0 348.6		057.7
Į.		- 1			1			.199	206.5	6.666	340.0	1.062	251.7
								.251	261.1	7.154	337-3	1.175	212.0
											-	S NA	





TABLE III.- MEASURED FLUTTER DERIVATIVES FOR THE NACA 65A004 AIRFOIL

a _m = 0°							α _m = 2 ⁰						
м	k	w	ga	φ	do	θ	ж	k	w	da da	φ	dc _m	8
-594	0.040	27.1	6.212	355.3	0.466	334-7	0.491	0.046	25.5	6.069	357.7		
ı	.080	53.4	5.849	350.1	-447	320.4		.095	52.5	5.699	2.6		
	.109	73.1	5.411 5.186	357.0 358.3	.517	308.1		.186	74.0	5.441	5.4 8.2		
	.188	125.7	5.195	8.2		300.1		.225	126.9	4.848	8.1		
	.224	153.3	4.883	353.5	.771	284.8		.267	150.3	4.367	12.4		
1	.261	178.0	5.242	353.0				.309	174.0	4.702	2.0		
	.342	233.5	6.989	353.8				.448	252.3	9.199	341.4		
ł	.382	260.9	8,995	330.7	1.396	219.8	.590	.040	26.9	5.862	355-7	0.656	343.0
.691	.035	26.9	6.846	356.1			.790	.080	53.0	6.011	357.3	.642	334.3
	.064	50.0	6.669	350.8				.112	74.3	5.873	359.8		
ŀ	.097	75.5	6.039	350.8 349.9				.154	102.4	5.362	1.2	.859	331.2
	.131	101.9	5.887	351.2				.193	131.4	5.091	351.8		
ĺ	.162	126.1	5.459	355.9 340.4				.226	154.0	4.982	351.6 347.7	.790	287.4
- 1	.194	153.9 204.9	5.251 4.981	3.0				.258 .338	231.0	5.501 6.340	4.1		
1	.258 .294	233.5	7.405	346.9				.370	252.3	7.528	2.7	1.282	256.7
- 1	.327	259.6	7.473	324.3		i				ļ.	1		
							.691	.033	25.6	7.400	355.6	.808	341.0
.741	.032	26.6	7.366	354.9	.561	338.2 327.8		.069	53.2 74.6	6.966	351.9 356.6	.839	330.2
1	.062	54.1 77.3	6.996 6.528	349.7 347.5	.578	321.0		.097	102.4	6.316	5.6	.856	332.8
	.117	101.9	6.235	345.9	.629	301.1		.164	126.9	6.397	356.4		
-	.145	126.5	5.861	339.7				.192	153.0	6.526	339.5	.984	298.9
ŀ	.169	145.3	5.078	329.0	.607	252.3		.285	227.1	8.342	350.1		
1	.239	205.3	5.599	357.6	.772	272.3		.321	255.8	9.530	330.7	1.564	243.7
- 1	.269	231.3	7.073	346.5	1.136	285.1	.741	020	CE 0	9 227	252.0	075	345.8
1	.303	260.7	7.035	337.8			. 141	.030	49.3	8.117 7.560	353.9 351.7	.975 .970	333.5
.798	.030	27.4	7.739	354.6	.609	340.3		.087	72.5	7.252	353.2	.510	333.7
~	.056	50.2	7.522	345.7	.658	320.5		.120	100.0	6.935	358.4	1.004	325.6
	.083	75.1	7.157	343.9				.147	122.7	6.574	359.1		
- 1	.115	103.7	6.606	343.5	.705	292.7		.175	151.6	6.128	332.8	1.034	276.5
]	.141	127.6	5.950	335.7				.235	202.7	5.990 8.895	354.9 339.4	.840	287.3
- 1	.197	183.7 203.0	5.301 5.477	343.5 343.2	.864	274.3		-297	256.8	8.784	333.0	1.536	244.6
	250	233.0	7.536	343.4	1.217	279.4		*=71	2,0.0	0.104	335.0	. ,,,,,	244.0
1	.278	259.1	7.924	323.1			.798	.029	26.0	9.165	352.9	1.140	344.5
.				•				.059	53-5	9.012	347.4	1.213	327.2
.850	.026	25.5	9.637	353.9				.083	74.5	8.584	351.7		318.2
	.053	50.9 74.8	8.200 7.611	344.4	.664	313.6		.114	102.5	7.322	349.1 346.1	1.209	310.2
Į	.106	102.2	7.025	337.4	.736	287.1		.191	178.8	6.260	349.0		
	.153	153.4	4,590	347.2	.588	276.5		.218	204.0	7.074	343.7	1.097	280.1
	.179	178.7	4.602	356.0				.244	228.5	9.002	339.2		
	.208	208.0	5.178	351.7	.928	274.2		.275	258.2	10.446	335.9	2.054	257.2
	.231	231.4	6.964	298.2	1.611	251.9	000	000	25.6		21.0 =	~~	221 77
	.259	258.5	8.436	325.4		:	.852	.060	25.6 52.3	9.672	348.5	.966 1.364	334.7 321.2
.900	.025	25.4	7.315	347.3	.717	331.8		.086	75.3	9.016	346.1 339.5		JE1.2
.,,,,	.050	51.1	8.913	344.6	.738	310.8		.117	103.0	8.075	345.2	1.227	301.6
	.073	74.1	8.483	337.4				.139	121.8	6.564	326.2		
- 1	.126	128.2	5.365	329.0				.176	177.5	5.269	345.7		
1	.147	155.4	4.721	341.9	.452	270.6		.204	205.1	5.928	353.7	1.430	282.7
- 1	.170	180.6	1.819	352.4	1.208	281.8		.228	230.2	8.365 9.206	336.3 338.2	2.450	279.5
	.195	206.6	8.097 8.076	0.6 337.7	1.947	305.2		1273		9.200	ع، در	2.+50	213.0
- 1	244	258.2	7.635	323.5	1.34	30).2	.870	.026	26.3	11.945	344.7		
		_						.049	49.7	11.356	330.1		
.942	.025	26.1	5.877	344.2	.438	328.4		.074	74.7	9.377	327.6	405	246.6
	.050	53.2	9.448	335.5	.824	290.3		.101	101.7	7.206	322.6	.510	229.7
	.096	102.2	6.885	316.1 334.7	.619	298.1		.142	174.9	4.826 5.713	342.6 356.1	.311	341.8
	.120	151.4	4.964	334.1				.190	202.7	6.013	350.5	.908	288.9
- 1		1						.212	226.4	7.880	340.0		
								.238	254.7	9.150	331.8	1.546	215.5
											ni= -		100 -
- 1							.904	.026	27.1	12.237 10.817	347.7 326.0	1.008	180.1 147.2
		ı l			1	1 1		.050		TO OT	J20.0	1,102	T+1.5
1		[1		.068	72.1	9.216	325.6		





TABLE IV. - MEASURED FLUTTER DERIVATIVES FOR THE NACA 2-008 AIRFOIL

			α <u>m</u> = 0°			-	α _m = 2°						
м	k	ω	da	φ	da da	е	к	k	G	da da	φ	de _m	6
0,590	0.040 .081 .113 .155 .193 .229 .350	26.3 53.9 75.3 103.0 128.7 152.5 232.7 259.6	6.460 6.082 5.624 5.436 5.425 5.287 5.001 6.210	354.6 353.5 351.4 352.3 351.2 351.2 7.0 1.2	0.552 .591 .620 .705	347.9 327.2 309.3 288.4 263.9	0.491	0.052 .093 .137 .181 .228 .278 .327 .465	28.6 51.5 75.9 100.2 126.2 154.0 181.1 257.5	6.240 5.919 5.838 5.329 5.324 5.543 5.396 5.638	358.2 355.0 351.4 355.5 356.2 1.2 3.1 22.5		
.680	.036 .069 .098 .134 .164 .197 .266 .299	27.9 53.0 75.5 103.9 127.2 152.4 206.0 231.0 254.7	7.006 6.433 6.118 5.600 5.437 5.227 4.627 5.723 5.931	351.6 347.9 349.0 347.1 346.0 343.6 358.3 356.1 350.6	.538 .615 .613 .692 .587	329.3 329.1 306.9 284.1 275.7	•590	.040 .080 .113 .146 .193 .231 .346 .384	26.8 53.9 75.6 98.2 129.6 155.1 231.9 257.5	6.691 6.204 6.093 5.706 5.650 5.665 5.379 6.678	354.2 346.8 348.4 352.9 349.6 348.5 4.8 1.4	0.574 .639 .604 .737	335.9 326.9 297.7 291.4 274.3
.728	.031 .063 .093 .123 .151 .249 .279	25.6 52.8 77.6 102.6 126.5 208.7 234.4 258.2	8.587 6.584 6.363 5.600 5.194 4.738 5.550 5.839	350.5 347.4 341.7 344.0 344.7 356.3 355.4 353.2	.788 .720 .738 .375	349.0 325.8 303.9 248.6 253.1	.680	.033 .069 .098 .134 .167 .200 .299	25.9 76.2 104.1 130.4 155.8 232.7 258.2	7.347 6.880 6.363 6.006 6.026 5.475 5.728 6.592	355.2 350.0 347.1 349.9 347.1 345.0 358.5 355.5	.661 .676 .645 .669	335.6 330.7 305.5 277.2 274.7
.786	.029 .056 .083 .111 .140 .196 .227	26.1 51.1 75.5 101.7 127.5 178.5 206.6 231.0	8.206 7.444 7.015 6.296 5.053 3.884 4.799 5.160	348.5 341.7 336.2 335.1 351.8 342.8 348.3	.665 .578 .659	338.4 315.3 285.2 256.9	.728	.031 .062 .090 .122 .152 .245 .279	26.0 75.7 75.7 103.0 127.9 206.7 235.6 258.2	7.829 7.517 6.957 6.932 6.624 6.264 7.365 7.482	358.2 354.5 353.5 349.3 344.6 2.6 358.2 358.9	.642 .441 .599	341.4 301.3 286.4 269.5
.833	.280 .028 .054 .080 .108 .164 .184	255.4 26.8 52.2 77.6 104.9 159.9 179.5	9.150 8.806 7.986 6.952 4.893 5.420	344.7 337.6 332.1 327.3 338.8 337.5	1.134 •399 •245 •357 •268	246.1 210.5 218.9 233.0 222.6	.786	.029 .056 .083 .113 .140 .227 .256	26.2 51.8 76.4 104.2 129.1 208.5 235.6 258.2	8.917 8.673 8.193 7.768 6.988 5.927 6.469 8.547	348.2 344.7 342.4 341.1 330.8 343.6 344.7 347.3	.086 .098 .260 .405	320.2 296.0 253.5 263.1 258.1
.879	.212 .239 .263 .027 .051 .076 .100 .150	206.5 232.7 256.1 27.6 53.1 78.2 103.0 155.3 177.5	5.576 6.807 7.713 9.720 8.316 7.062 5.235 3.831 4.760	346.4 348.3 337.3 348.8 332.1 325.5 324.2 340.4 342.5	.538 .525 1.250 1.353 .986 .454	227.9 214.1 156.3 138.2 139.5 135.3	.833	.028 .053 .078 .107 .185 .212 .239	27.1 52.4 76.9 104.7 181.8 207.6 234.2 257.2	9.455 9.008 8.618 7.950 5.425 5.662 7.075 8.568	353.9 345.4 338.8 327.7 342.0 349.8 349.6 341.2	.317 .316 .339 .346 .467	203.3 200.2 172.2 261.4 281.4
	.200 .226 .252	206.0 233.3 259.6	6.218 6.432 7.107	341.4 335.2 330.6	.642 1.007	165.4 157.7	.879	.027 .052 .075 .100 .148 .174 .201 .225 .248	28.7 53.8 78.5 104.3 154.2 181.6 209.4 234.4 258.6	10.660 8.749 6.975 5.550 4.421 4.481 6.225 6.098 6.522	343.1 332.4 326.0 326.9 335.5 357.0 348.7 342.3 331.1	1.437 .527 .485 .743	142.5 144.1 96.8 179.0

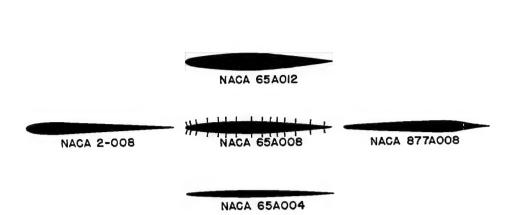




TABLE V.- MEASURED FLUTTER DERIVATIVES FOR THE NACA 877A008 AIRFOIL

		,	°= 0		,		<u> </u>			c _{11 = 2} <			
ĸ	k		do	•	do.	•	ж	k		de	P	don	
0.495	0.050	27.9	6.087	352.7			0.496	0.048	27.1	6.416	353.0		
	.092	51.5 75.7	5.638 5.136	349.5	1:::		1	.087	49.1	5.936	350.8		
	.183	103.0	5.159	348.2				.128	72.1	5.497	351.9 349.7		
	.228	128.2	4.941	355.0				.178	125.4	5.429	347.8		
	.266	149.6	4.956	350.9				269	151.4	1.985	347.8		
	.316 .450	179.5 255.4	5-135	338-1			1	.321 .454	179.5	5.243	352.8	~	
	1	l	1	4.8	1			.454	253.3	4.369	18.3		
-596	.040	27.2 49.8 74.3	6.246	348.0	.438	339.8	.596	.035	24.0 51.3	6.579	1.7 350.6	0.634	342. 333.
	.109	74.3	5.462	346.5				.108	73.6	5.948	348.7		333.
	-153	103.7	15.448	345.2	-609	310.6		.147	99.6	5.612	350.8	.765	320.
	.187	126.9 151.7	5.370 4.728	345.5 346.0	F60	202.0		.182	123.9	5.798	352.7		
	-336	231.0	4.267	355.1	-560	301.9		.218	148.2	5.067	347.1	.625	213.
	-375	257.5	5.053	3.4	.963	293.8		.259 .337	227.6	5.325 4.688	340.3 357.6		= =
.693	.035	27.5	6.976	348.5				.372	251.3	5.319	3.6	1.192	197.
	.066	76.7	6.842	343.4			.693	.031	25.1	6.333	354.2	.536	٥.
	.096	102.5	6.423	343.0		~		-064	51.3	6.034	352.4	-551	334.
	.158	125.8	5.629	339.5				.092	73-9	5.614	346.0		
	.192	152.7	5.002	336.3				.127	101.4	5.562	345.7	-637	310.
	-254	204.4	3.861	0.0				.153	122.7	5.202	342.1 335.9	.659	287.
- 1	-285	229.9	4.503	356.8				257	201.2	3.129	334.9	1.487	192.
	.321	258.9	4.054	356.1				-257 -289	229.3	016	0.1		
.745	.031	26.7	7.230	349.8	.705	336.0		.319	253.3	3.884	358.0	1.044	176.
	.060	52.1	6.933	346.6	.664	332.7	.745	.030	26.3	7.318	351.3	.744	352.
j	-086	74.1	6.693	341.4				.058	50.5	7.049	349.0	.819	337.
	.118	102.0	6.192	341.1	-743	310.8		.088	50.5 75.8	6.674	343.7		
1	.236	206.0	5.863	332.1 354.9	.613	312.6		.116	99.8	6.328	342.7	1.006	314.
	.262	228.4	4.719	345.5	.013	312.0		.146	126.0	5.912	337 - 3 355 - 4		
	.291	253.3	1.218	343.6	.906	291.7		.237	204.0 235.6	4.875 5.120	337.9	.832	214.
.796	.029							-296	255.1	4.560	345.7	.838	210.
.190	.056	26.9	8.056 7.454	352.9 345.6	.881 .599	345.6	.798	.028	26.1	9 000	354.4	-777	1.
- 1	-081	75.5	7.497	337-9			3,,50	054	50.6	7.410	340.8	.791	322.
- 1	.110	102.2	6.566	334-3	.765	303.4		079	73.6	7.030	339.8		
	.137	127.3	3.642	330.1 349.5	===		1	.109	101.3	6.380	331.4	.804	294.
		207.3	4.442	348.0	بانان	317.3		.135	125.7	5.297	322.9		~ -
	.250	234.0	4.766	343.6		251.2		.190	176.5	4.123	335.8 345.3	.519	208.
	.280	262.5	5.340	347.5	1.177	289.5		.221	205.3		325.7	.519	200.
SOF	~~~			-10 -				.279	259.6		339.8	.830	228.
.825	.029	28.2 5 2.1	7.268	348.2	1.168	331.0	.827						
- 1	.078	75.6	6.764	340.2	1.113	325.6	.027	.027	25.9	8.460	346.8		
- [.106	102.7	6.302	331.3	-997	308.1	- 1	.053	51.5 76.1	7.945	344.1		
Ī	-130	125.8	5.291	330.6			1	.107	103.4	6.133	330.2		
ł		178.5	1.023	356.1			1	185	178.5	6.133 2.884	345-7		
- 1	.213 .238	207.3	4.581	35.5	1.389	329.1		.210	203.3	3.100 4.054	2.7		
		257.5	5.601	340.9	1.490	289.5	1	.240	231.8	4.054	347.1		
_			Ĭ		. 1			.255	246.4	4.952	348.9		
-857	027	27.3 51.1	6.705	351.3 345.8	1.637	349.2 338.9	.860	.026	26.3 51.6	8.294 7.660	351.4 334.9	.997 .761	355. 320.
	.074	75.1	6.048	336.9		350.9	1	-074	71.3	5.404	324.6	.101	
	-103	103.9	5.451	325.8	1.342	320.1		.098	99.2	5.707	315.4	.776	283.
- 1	.152	153.2	3.359	342.3	.722	327.6	1	,153	154.2	3.976	331.0	.643	298.
- 1	.202	181.6	2.715	338.3	1 519	350 h	1	.178	179.5	4.292	326.1	770	183.
1	229			333.6	1.518	352.4	i	.204	205.8	4.064	335.9	.770	105.
	253	256.5	3.282	334.3	2.161	307.4		-252	254.7	3-773	317.3	1,606	148.
.883	.027	27.8	6.494	350.5	1.702	0.6	.892	.025	26.1	9.005	340.4	.446	٥.
	.051	53.1	6.162	338.1	1.899	328.1	1	.049	50.3	7.821	333-4	1.172	209.
- 1	.100	76.0	5.577 4.493	332.4 325.7	1.447	322.8	1	.073	75.3 151.7	3.753	317.0 334.9	.520	327.
- 1	.146	12.1	3.493	336.2		335.1	1	.173	161.6	3.123	329.0		3= 1 ·
- 1	.171	179.0	3.023	336.9			Ė	.202	212.2	3.685	332.2 323.5	.172	171.
	222	232.7	3.928	338.5 338.7	1.951	309.5		.224	212.2 235.3 257.5	3.168	323.5 314.0	.755	322.
مدو.						,.,			,,-,			-,,,,	
	-135	145.8	3.920	343.7 343.3		162.6							
- 1		179.5	4.358	342.5									
- 1	.197	230.6	1.021			174.9							
- 1				299.7		173.0	i						
						-12.0			. 1		1		

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MODEL PRESSURE-CELL LOCATIONS
[In Percent of Model Chord]

Cell number upper and lower surface	65A012 and 65A008	65A004 2-008, and 877A008
1 2 3 4 5 6 7 8 9 0 1 2 3 4 5 15	1.25 3.75 7.5 15 22.5 27.5 35 45 57.5 67.5 67.5 75 85 95	1.25 3.75 7.5 15 22.5 27.5 35 45 57.5 67.5 67.5 85 90



Figure 1.- Section profiles and pressure-cell locations of models.

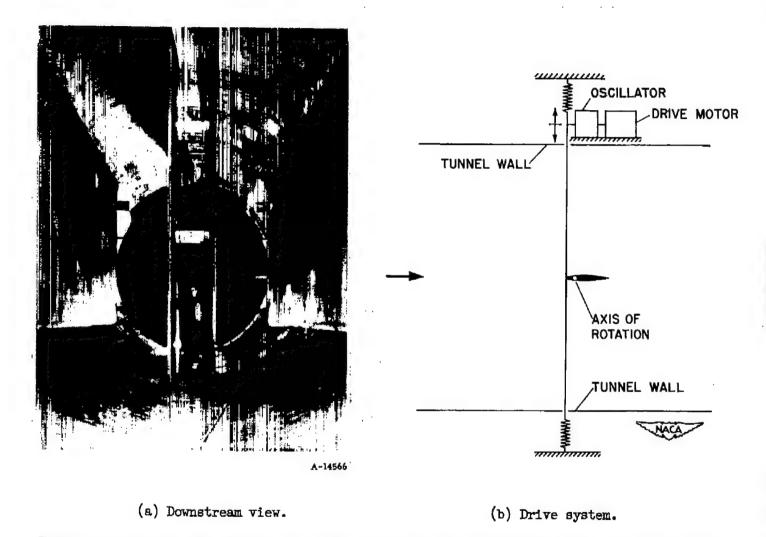


Figure 2.- View of test section with model in place and diagrammatic sketch of drive system.

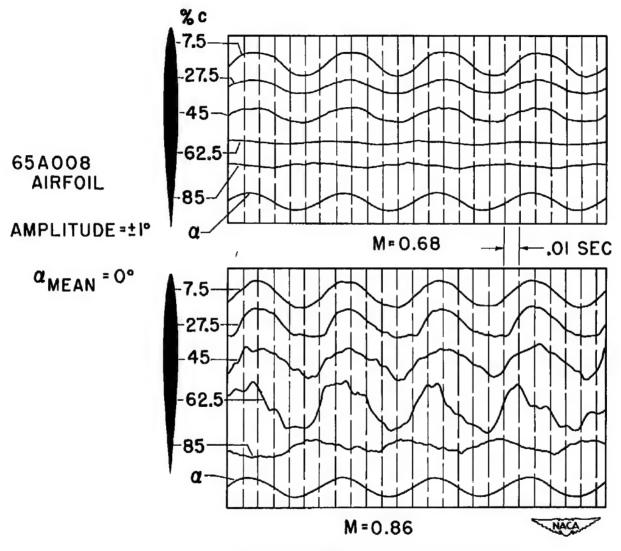


Figure 3.- Typical oscillograph traces.

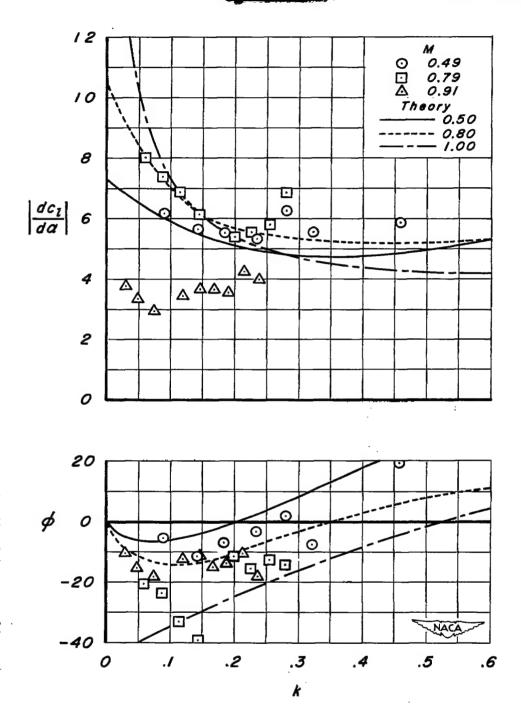
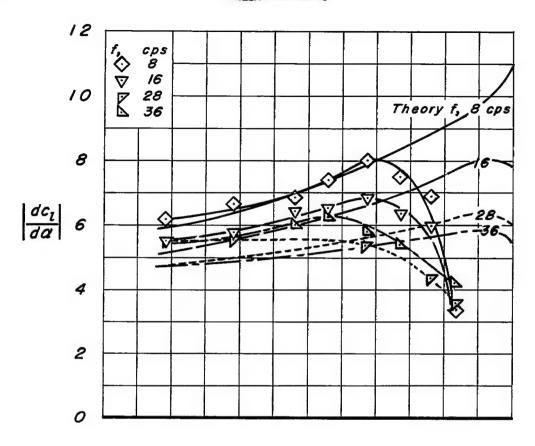


Figure 4.— Results as a function of reduced frequency, k, for several Mach numbers for the reference model, NACA 65A008; $\alpha_m = 0^{\circ}$.



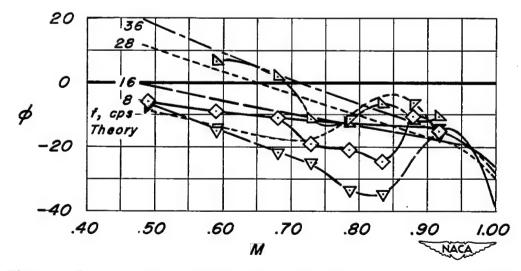
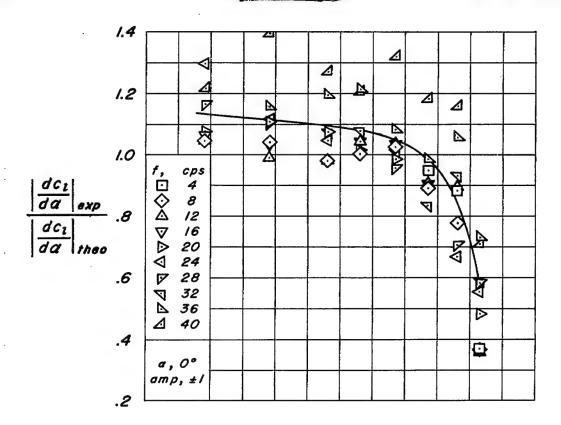


Figure 5.- Typical results for reference model, NACA 65A008; $\alpha_m = 0^\circ$.





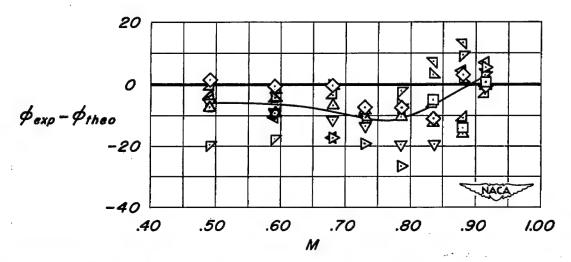
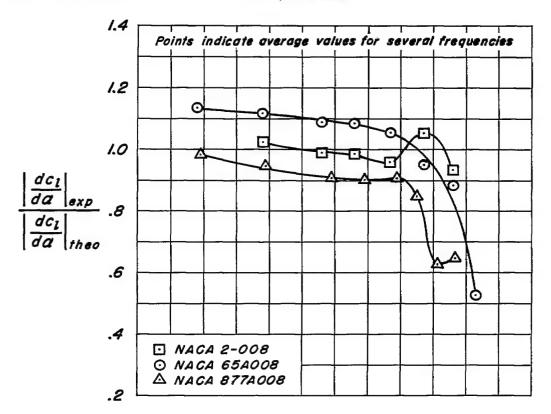


Figure 6.- Variation of experimental results from theory for reference model, NACA 65A008, with a faired line to show the mean variation with Mach number; $\alpha_m = 0^{\circ}$.







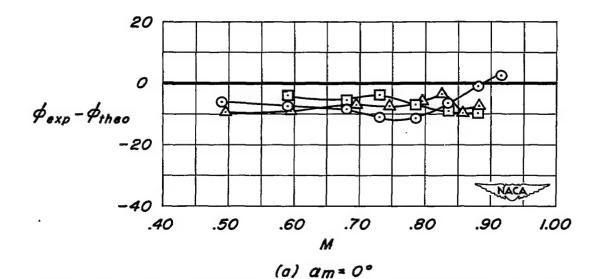
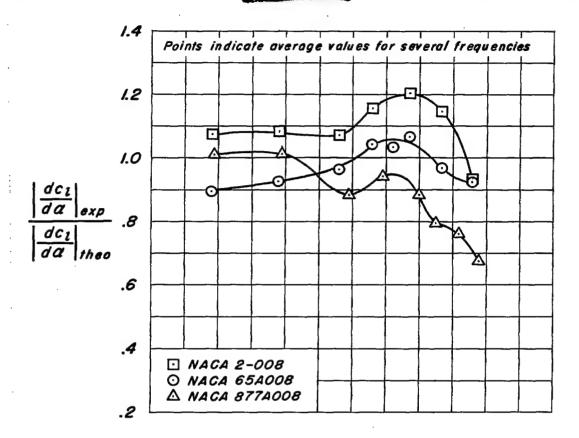
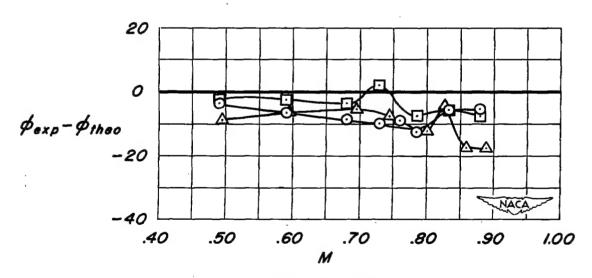


Figure 7.- Effect of airfoil thickness distribution on lift derivatives.

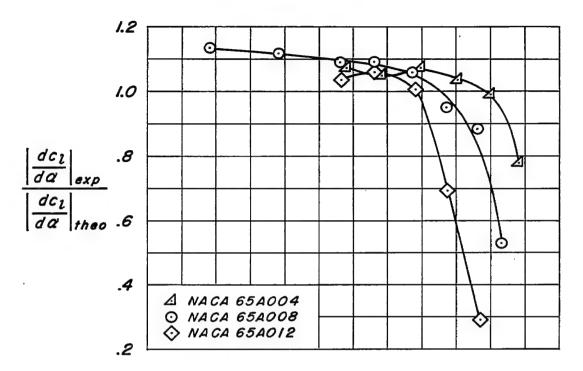






(b) a_m = 2° Figure 7.- Concluded.





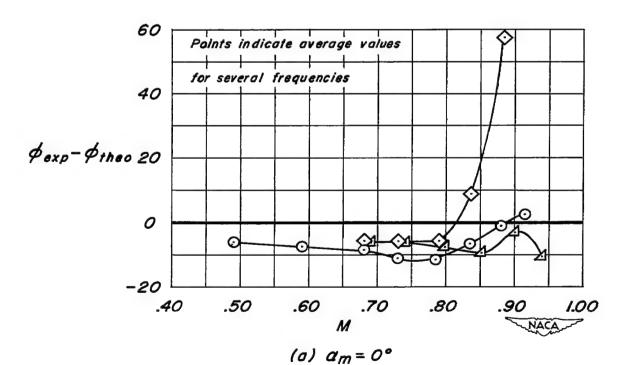
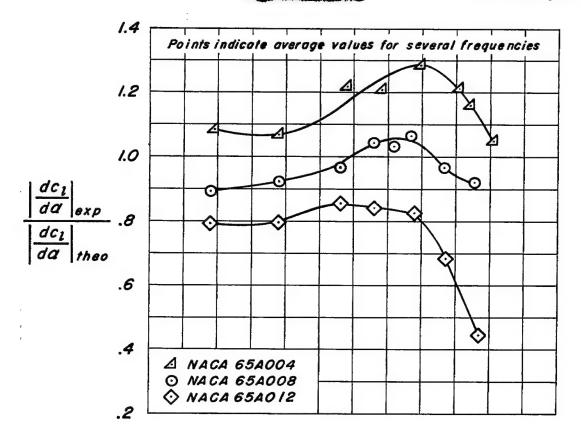
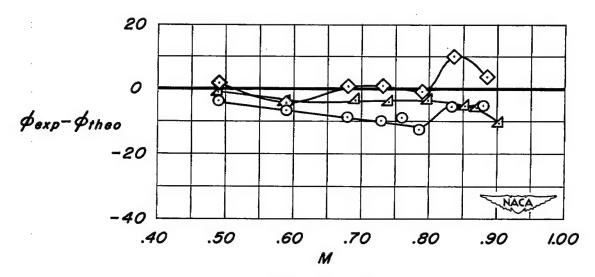


Figure 8.- Effect of airfoil thickness on lift derivatives.







(b) $a_m = 2^{\circ}$ Figure 8.- Concluded.

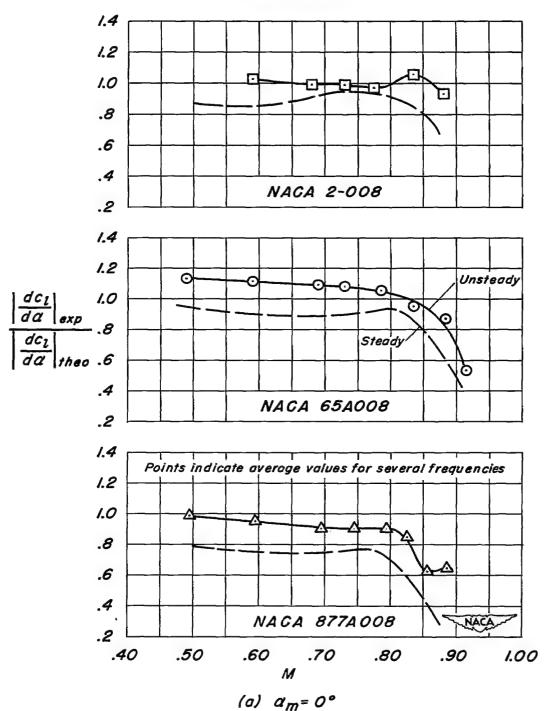
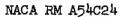
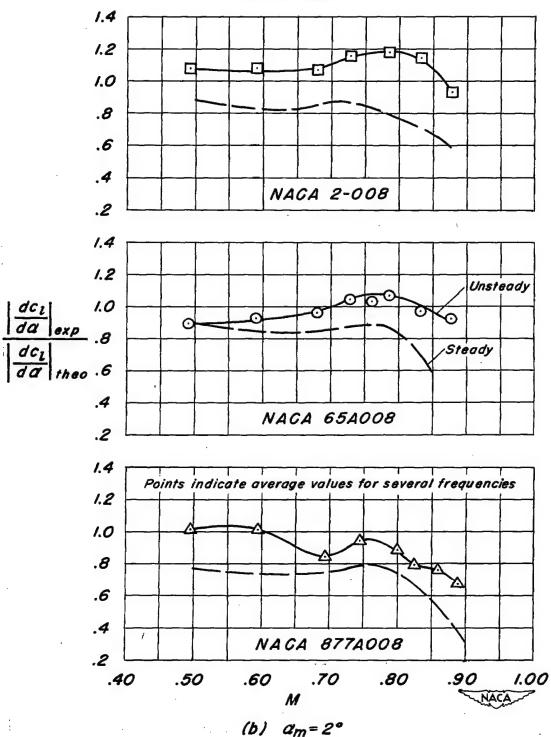


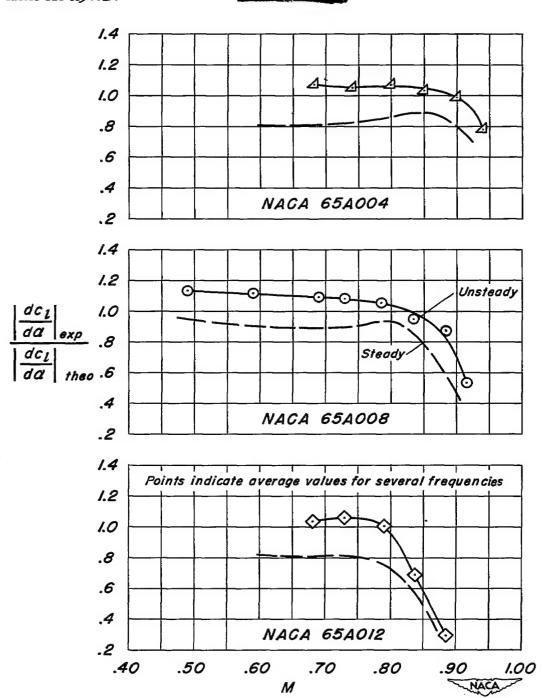
Figure 9.- Comparison of steady and unsteady lift derivatives for airfoils with varying thickness distributions.



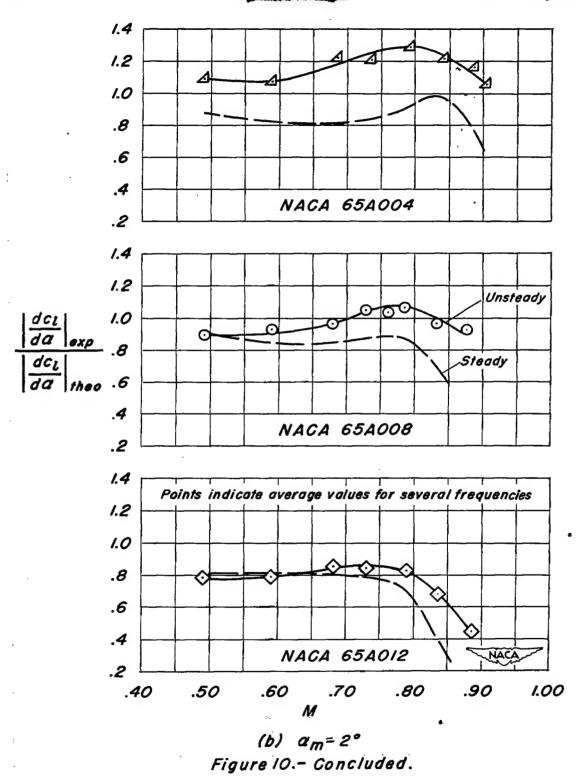


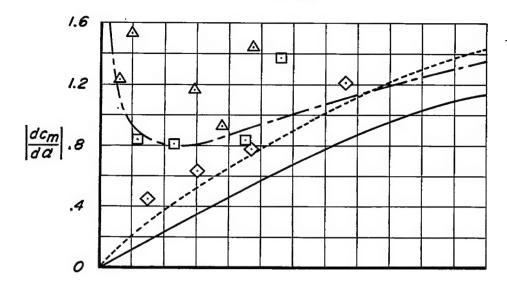
(D) $\alpha_m = 2^{\circ}$ Figure 9- Concluded.

 $\sigma^* \in \mathbb{R}$



(a) $\alpha_m = 0^\circ$ Figure 10:— Comparison of steady and unsteady lift derivatives
for airfoils with varying thickness.





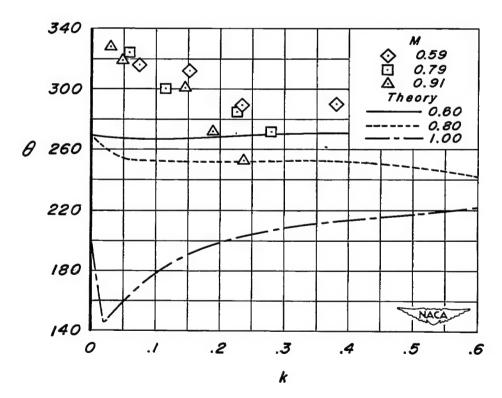


Figure II.- Results as a function of reduced frequency, k, for several Mach numbers for the reference model, NACA 65A008; $\alpha_m = 0^\circ$.



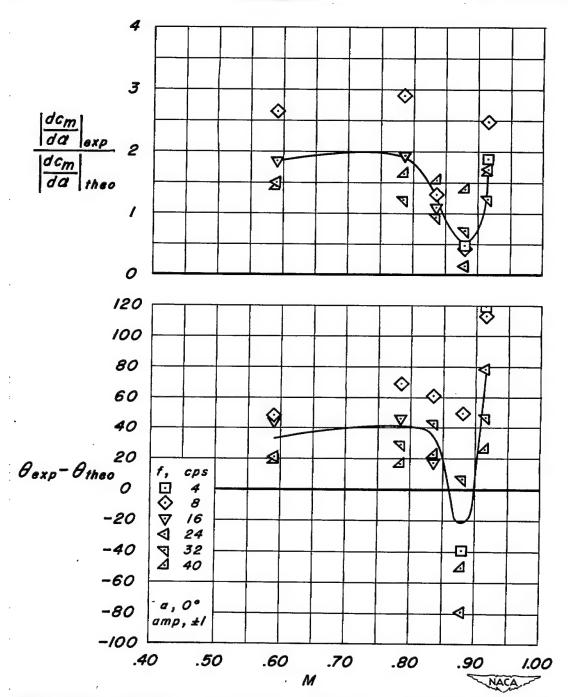
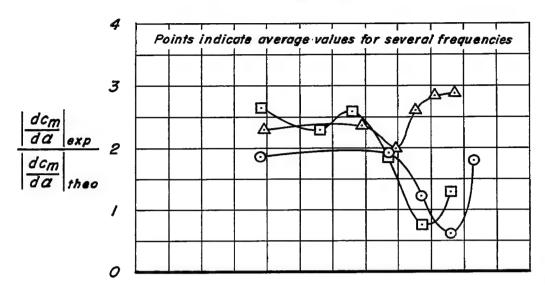


Figure 12.— Variation of experimental results from theory for reference model, NACA 65A008, with a faired line to show the mean variation with Mach number; $a_m = 0^\circ$.





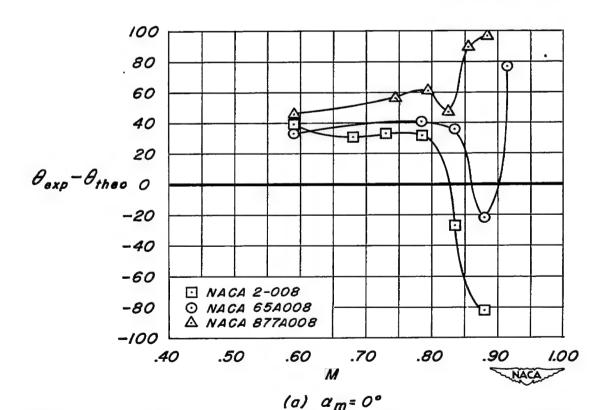
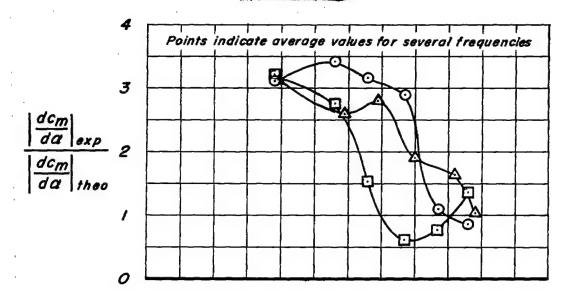
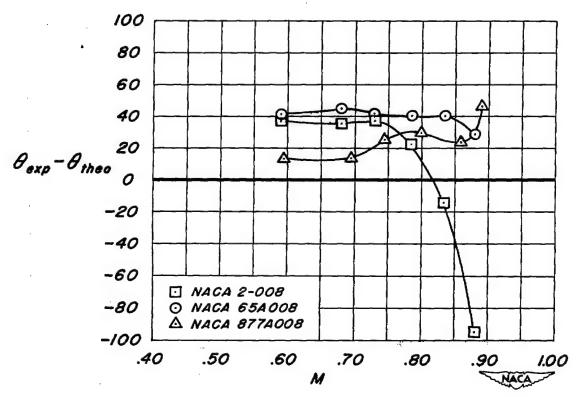


Figure 13.- Effect of airfoil thickness distribution on moment derivatives.







(b) $a_m = 2^{\circ}$ Figure 13.- Concluded.



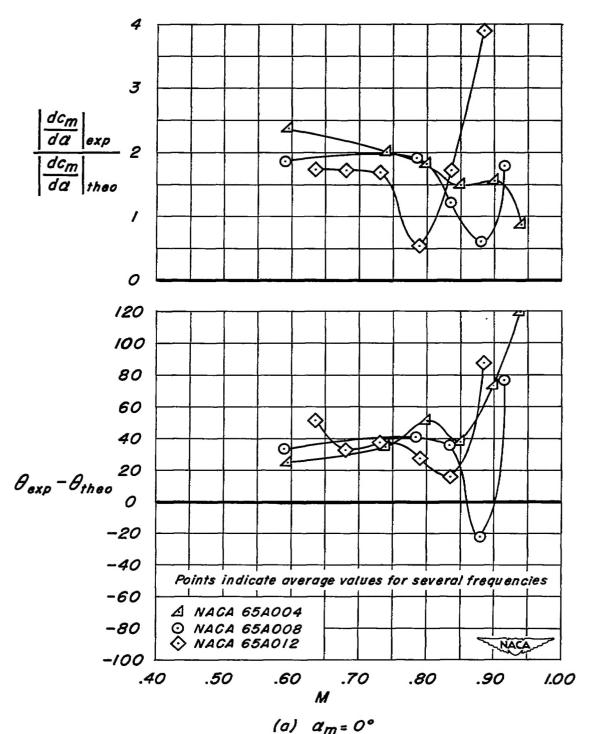
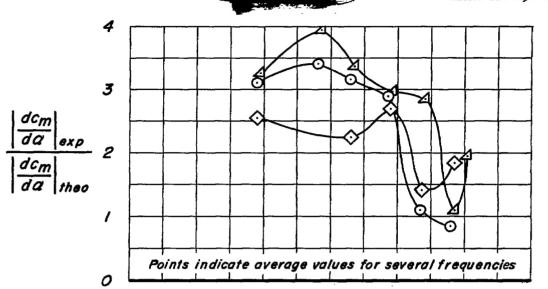
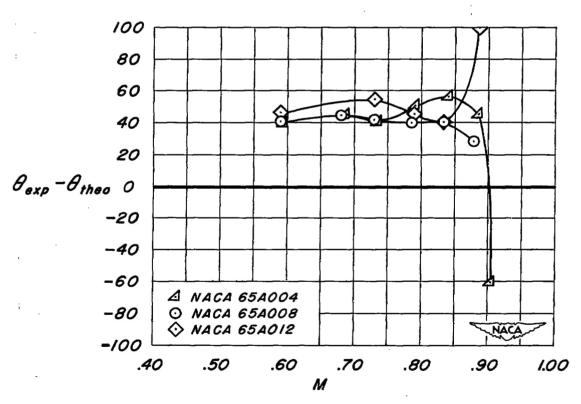


Figure 14.- Effect of airfoil thickness on moment derivatives.

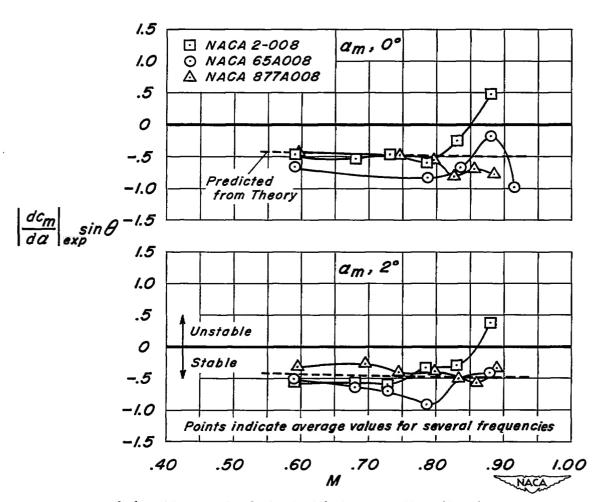




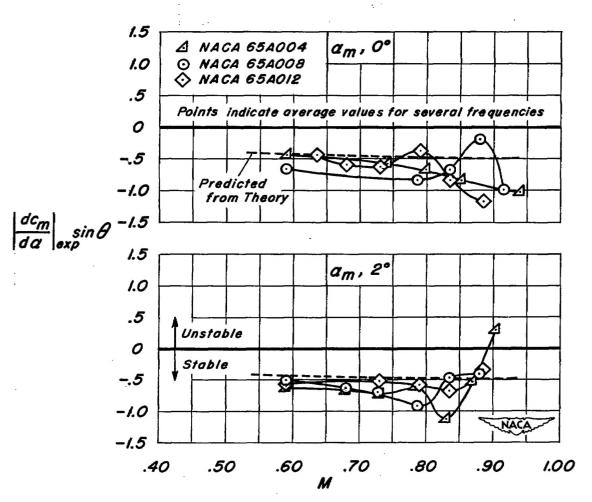


(b) $a_m = 2^\circ$ Figure 14.- Concluded.





(a) Effect of airful thickness distribution.
Figure 15.— Damping component of the moment derivatives.



(b) Effect of airfoil thickness. Figure 15.- Concluded.